
Chameleons in the Laboratory

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Outline

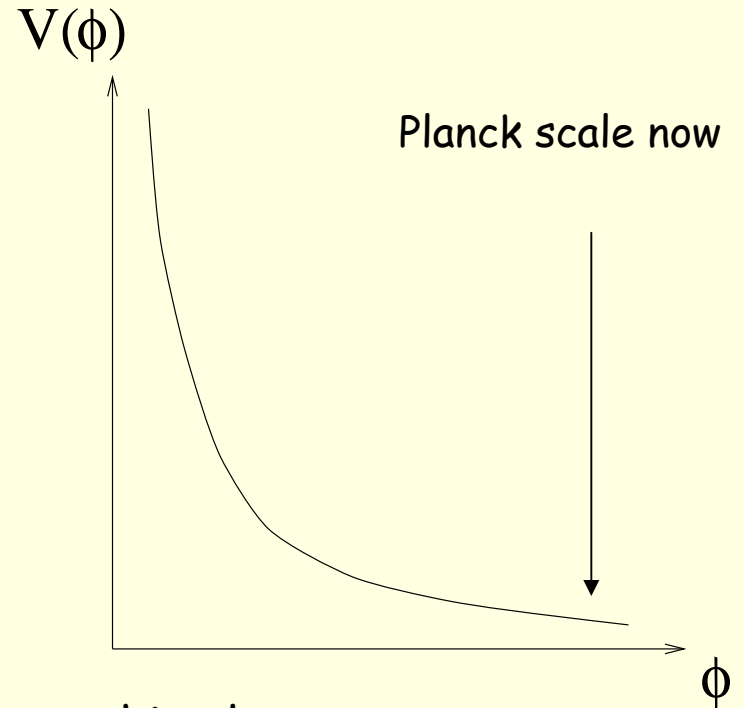
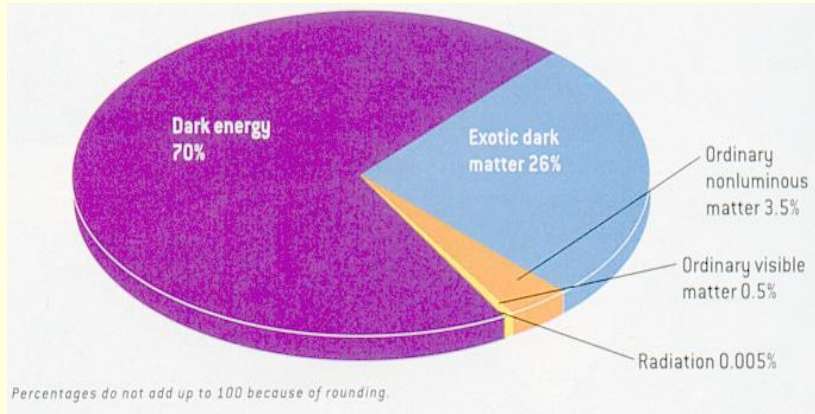
1-Why Chameleons? (Link with Dark Energy)

2- Chameleon Optics

3- Chameleonic Casimir Effect

Why Chameleons?

Dark Energy



Field rolling down a runaway potential, reaching large values now (Planck scale)

Extremely flat potential for an almost decoupled field

How Flat?

Energy density and pressure:

$$\rho = \frac{1}{2}\dot{\phi}^2 + V(\phi), \quad p = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$

Runaway fields can be classified according to

$$w = \frac{p}{\rho}$$

$$m \gg H_0$$

very fast roll

$$w \approx 1$$

$$m \ll H_0$$

slow roll

$$w \approx -1 \text{ (inflation)}$$

$$m \approx H_0$$

gentle roll

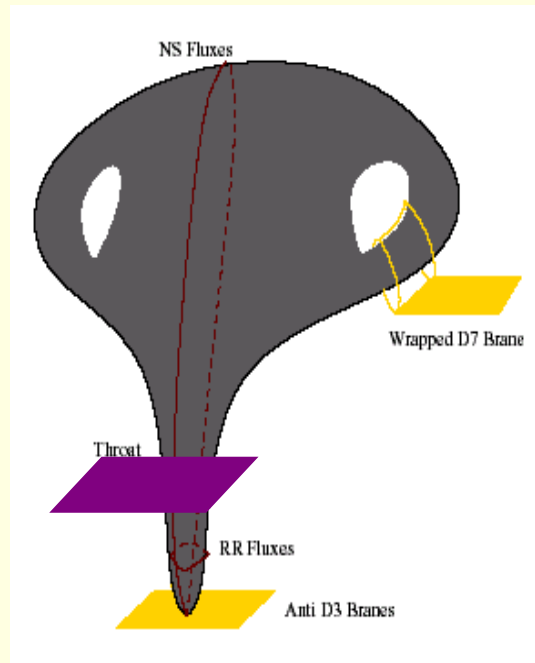
$$w \neq -1 \text{ (quintessence)}$$

$$H_0 \approx 10^{-43} \text{ GeV}$$

→ strong gravitational constraints

Moduli

KKLT scenario



To all orders of perturbation theory, the moduli have a flat potential. Non-perturbative effects can lead to runaway potentials.

Beyond the standard model physics leads to the existence of myriads of scalar fields: the moduli.

Moduli Couple to Matter!

The Standard Model fermion masses become moduli dependent

$$m_{u,d}(\phi) = \lambda e^{\kappa_4^2 K(\phi)/2} m_{u,d}$$

Scalar-tensor theory

Yukawa

Kahler

$$K = -\frac{n}{\kappa_4^2} \ln \kappa_4(\phi + \bar{\phi})$$

$n=1$ dilaton, $n=3$ volume modulus

Gravitational Problems

- o Deviations from Newton's law are tested on macroscopic objects. The gravitational coupling is:

$$\kappa_4 \alpha = \frac{d \ln m_{\text{atom}}}{d\phi_n} \longleftarrow d\phi_n = \sqrt{K_{\phi\phi}} d\phi$$

- o The deviation is essentially given by:

$$\alpha \approx \left[\frac{3m_u + m_d}{2 \Lambda_{\text{QCD}}} - \frac{1}{8} \frac{N - Z}{N + Z} \frac{m_u - m_d}{\Lambda_{\text{QCD}}} \right] \kappa_4 \partial_{\phi_n} K$$

- o For moduli fields:

$$\kappa_4 \partial_{\phi_n} K = \sqrt{2n}$$

Too
Large!

Gravitational Tests

Scalar-tensor theories suffer from the potential presence of a fifth force mediated by the scalar field.

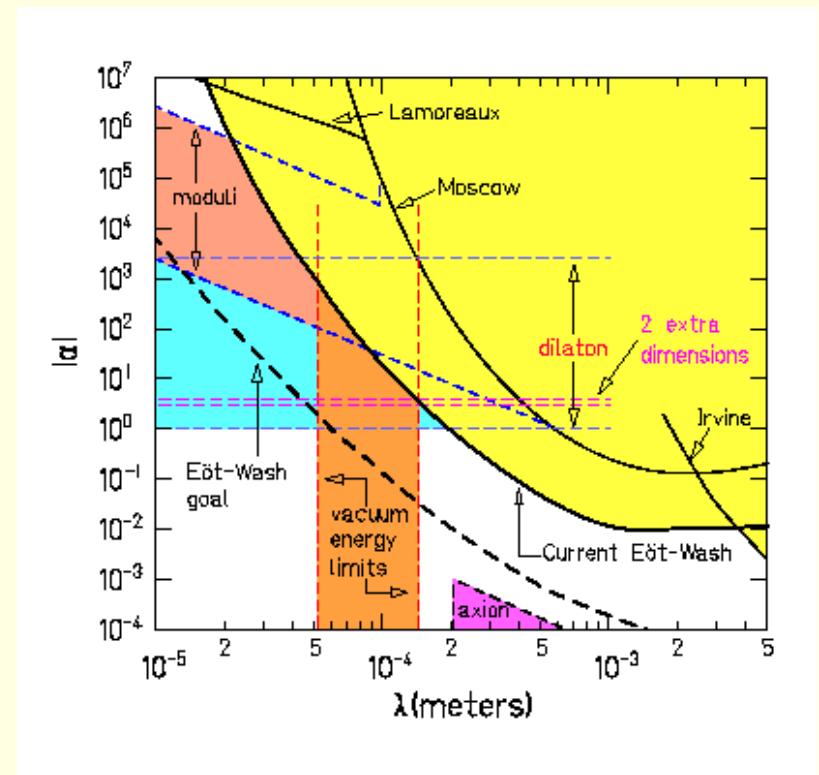
Alternatives:

Non-existent if the scalar field has a mass greater than :

$$m \geq 10^{-3} \text{ eV}$$

If not, strong bound from Cassini experiments on the gravitational coupling:

$$\alpha^2 < 10^{-5}$$



Chameleons!

Chameleon field: field with a matter dependent mass

A way to reconcile **gravity tests and cosmology:**

Nearly massless field on cosmological scales

Massive field in the laboratory



Chameleon Effective Theory

Effective field theories with gravity and scalars

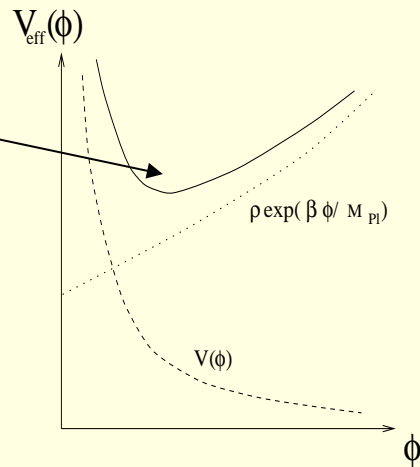
$$S = \int d^4x \sqrt{-g} \left(\frac{1}{16\pi G_N} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) + \mathcal{L}_m(\psi_m, A^2(\phi)g_{\mu\nu}) \right)$$

The Chameleon Mechanism

When coupled to matter, scalar fields have a matter dependent effective potential

$$V_{\text{eff}}(\phi) = V(\phi) + \rho_m A(\phi)$$

Environment
dependent
minimum



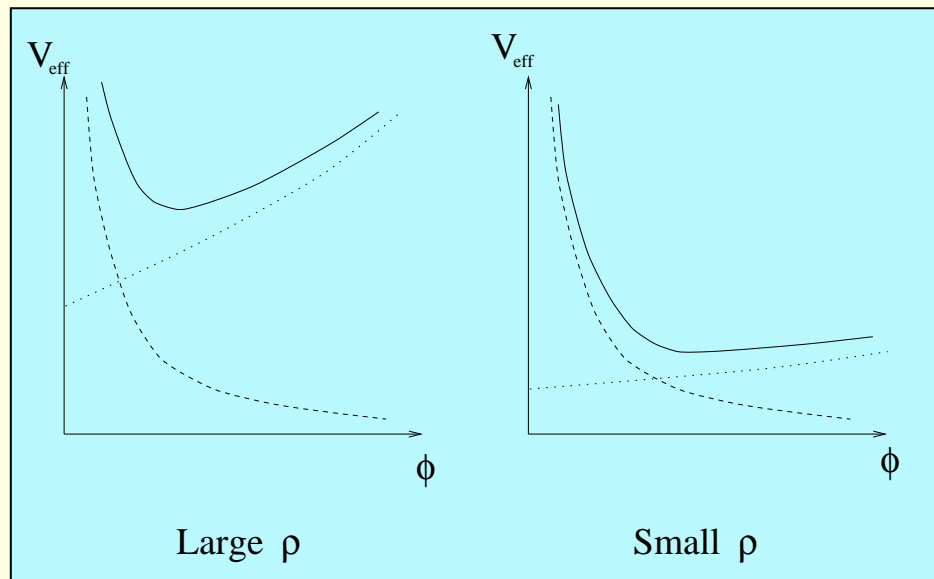
An Example:

Ratra-Peebles potential

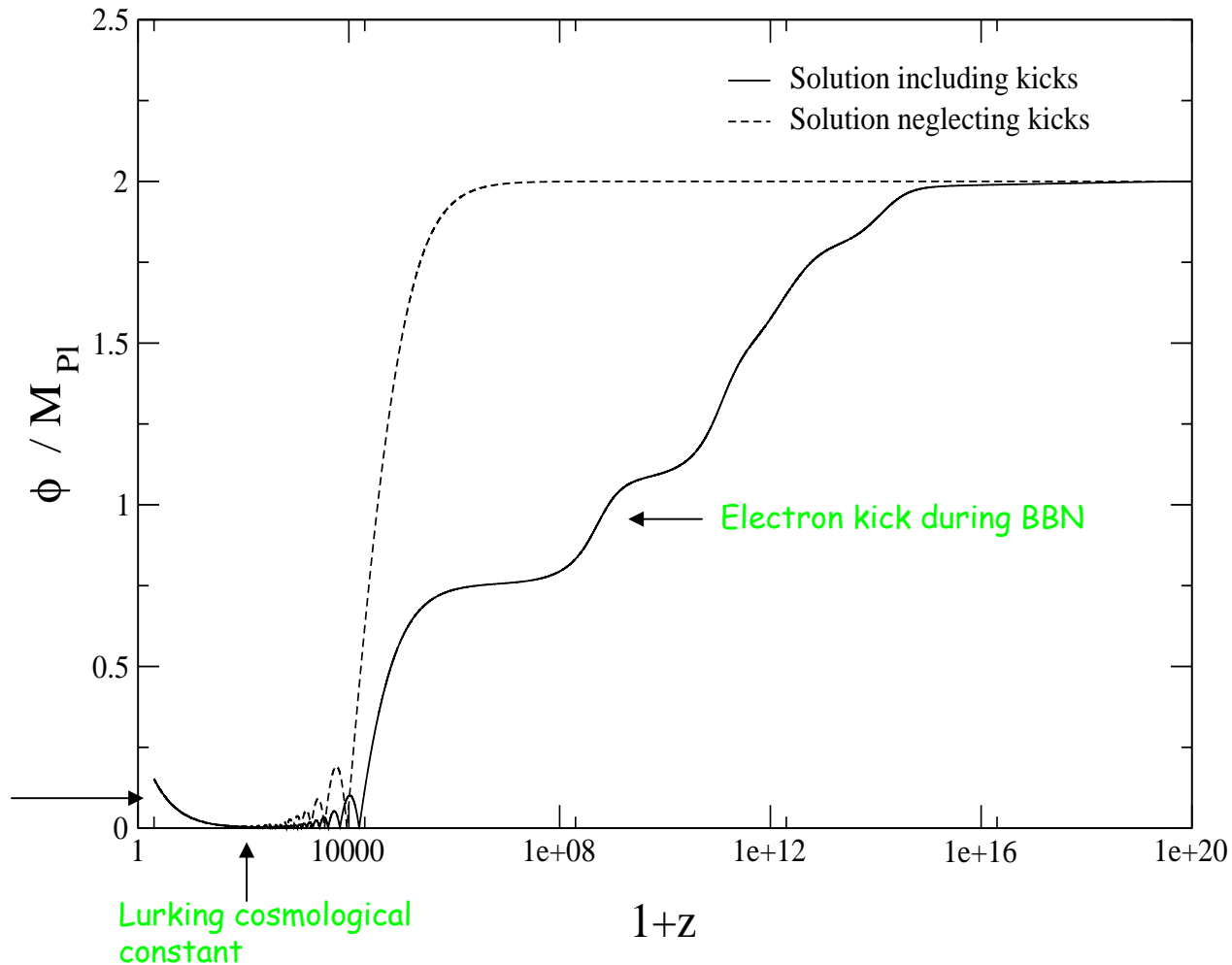
$$V(\phi) \sim \frac{\Lambda^{4+n}}{\phi^n}$$

Constant coupling to matter

$$A(\phi) = \exp \frac{\phi}{M}$$



Chameleon Cosmology



Chameleon Optics

Chameleons Coupled to Photons

- Chameleons may couple to electromagnetism:

$$\mathcal{L}_{\text{optics}} = \frac{e\phi/M}{g^2} F_{\mu\nu} F^{\mu\nu}$$

- Cavity experiments in the presence of a constant magnetic field may reveal the existence of chameleons. The chameleon mixes with the polarisation orthogonal to the magnetic field and oscillations occur (like neutrino oscillations)

- The coherence length $z_{\text{coh}} = \frac{2\omega}{m^2}$

depends on the mass in the optical cavity and therefore becomes pressure and magnetic field dependent:

$$\rho = \rho_m + \frac{B^2}{2}$$

- The mixing angle between chameleons and photons is:

$$\theta = \frac{B\omega}{Mm^2}$$

Ellipticity and Rotation

- Photons remain N passes in the cavity. The perpendicular photon polarisation after N passes and taking into account the chameleon mixing becomes:

$$\psi(z) = N \left(1 - \frac{1}{N} \sum_{n=0}^{N-1} a_n(z) \right) \cos \left(\omega z + \frac{1}{N} \sum_{n=0}^{N-1} \delta_n(z) \right)$$

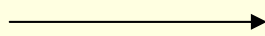
- The phase shifts and attenuations are given by:

$$a_n(z) = 2\theta^2 \sin^2 \frac{m^2(z+nL)}{4\omega}, \quad \delta_n(z) = \frac{m^2\theta^2}{2\omega}(z+nL) - \theta^2 \sin \frac{m^2(z+nL)}{2\omega}$$

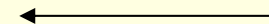
identified with the phase shift and attenuation after one pass of length nL.

- At the end of the cavity $z=L$, this can be easily identified for commensurate cavities whose lengths corresponds to P coherence lengths

Rotation



$$a_T = \theta^2, \quad \delta_T = \pi \frac{N}{P} \theta^2$$



ellipticity

Ellipticity vs Rotation

- o For a large number of passes, the ellipticity is always larger than the rotation

$$\frac{\text{ellipticity}}{\text{rotation}} = \pi \frac{N}{P}$$

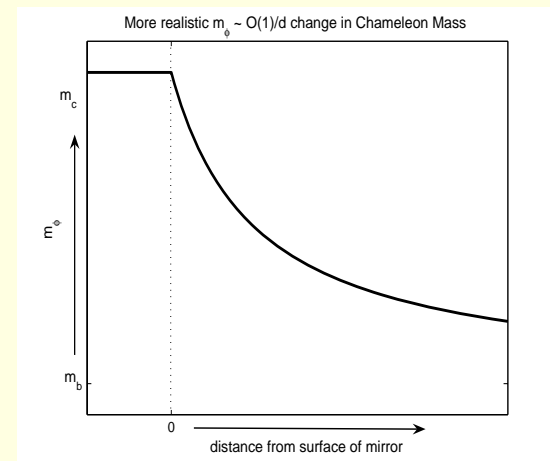
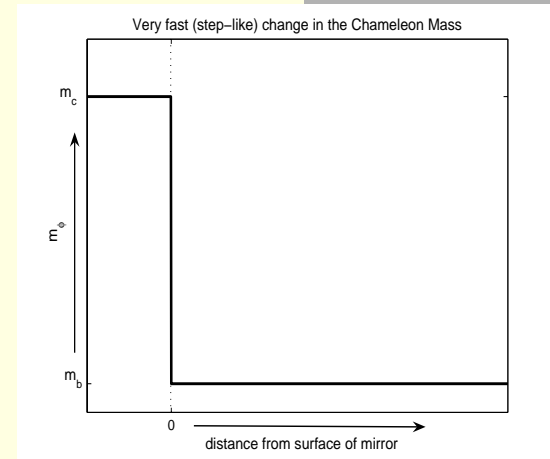
Realistic Chameleon Optics

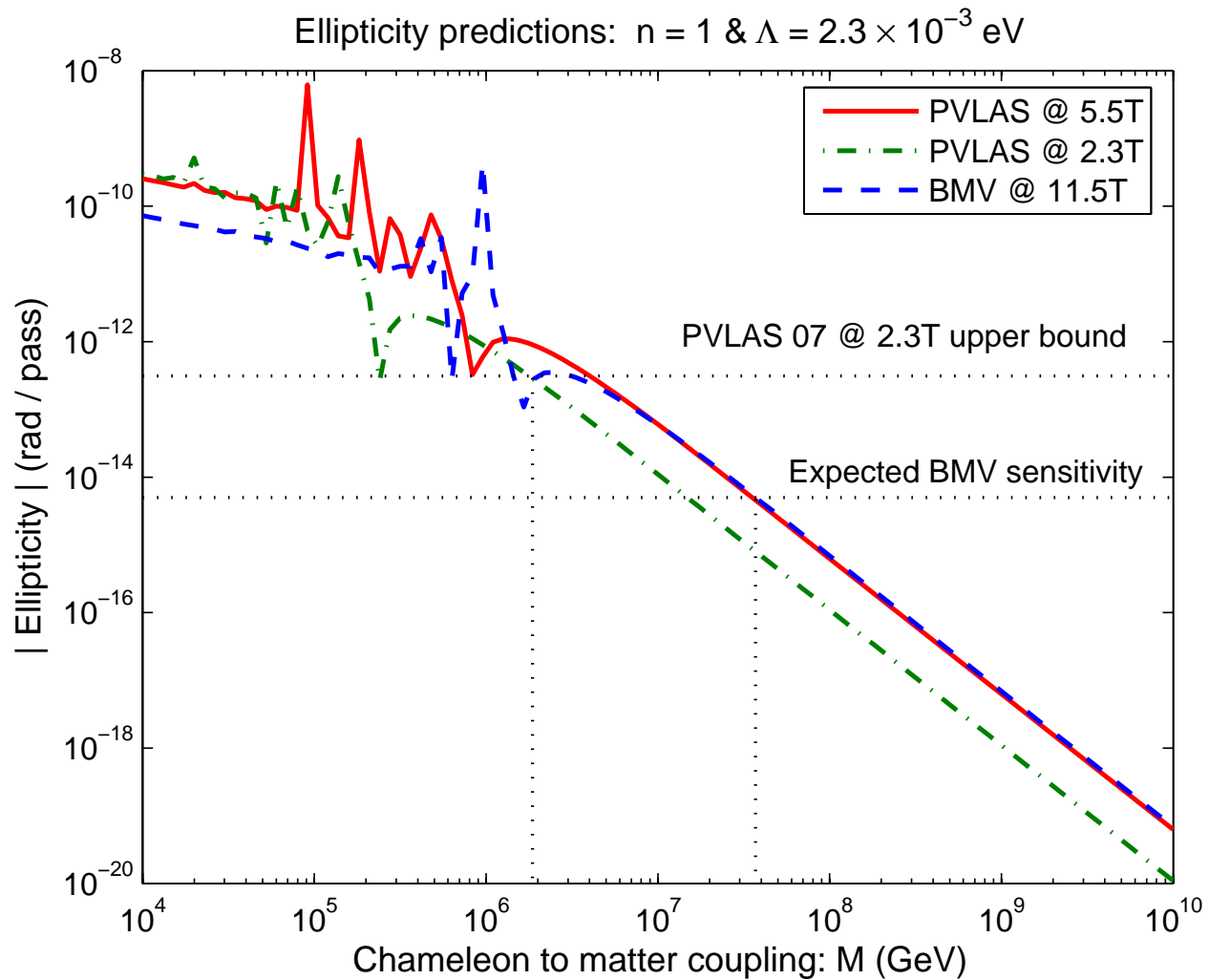
- Must take other effects into account.
- Chameleons never leave the cavity (outside mass too large, no tunnelling)
- Chameleons do not reflect simultaneously with photons.

$$\Delta_r = \frac{\pi n}{n + 2}$$

- Chameleons propagate slower in the no-field zone within the cavity

$$\Delta_d = \frac{m_\phi^2 d}{\omega}$$





Chameleonic Casimir Force

The Thin Shell Effect I

- o The force mediated by the chameleon is:

$$F_\phi = -\beta \frac{m}{m_{\text{Pl}}} \nabla \phi, \quad \beta = \frac{m_{\text{Pl}}}{M}$$

- o The force due to a compact body of radius R is generated by the gradient of the chameleon field outside the body.
- o The field outside a compact body of radius R interpolates between the minimum inside and outside the body
- o Inside the solution is nearly constant up to the boundary of the object and jumps over a thin shell ΔR
- o Outside the field is given by:

$$\phi \approx \phi_\infty - \frac{\beta}{m_p} \frac{3\Delta R M_c}{R r}$$

The Thin Shell Effect II

- The force on a test particle outside the body is proportional to Newton's attraction

$$F_\phi = 6\beta^2 \frac{\Delta R}{R} F_N$$

- When the shell is thin, the deviation from Newtonian gravity is small.
- The size of the thin-shell is:

$$\frac{\Delta R}{R} = \frac{\phi_\infty - \phi_c}{6\beta m_p \Phi_N}$$

- For bodies used in the Casimir experiments, this is always the case when β is greater than 1000.

The Casimir Force

- o We focus on the plate-plate interaction in the range:

Mass in the plates \longrightarrow $m_c^{-1} \leq d \leq m_b^{-1}$ \longleftarrow Mass in the cavity

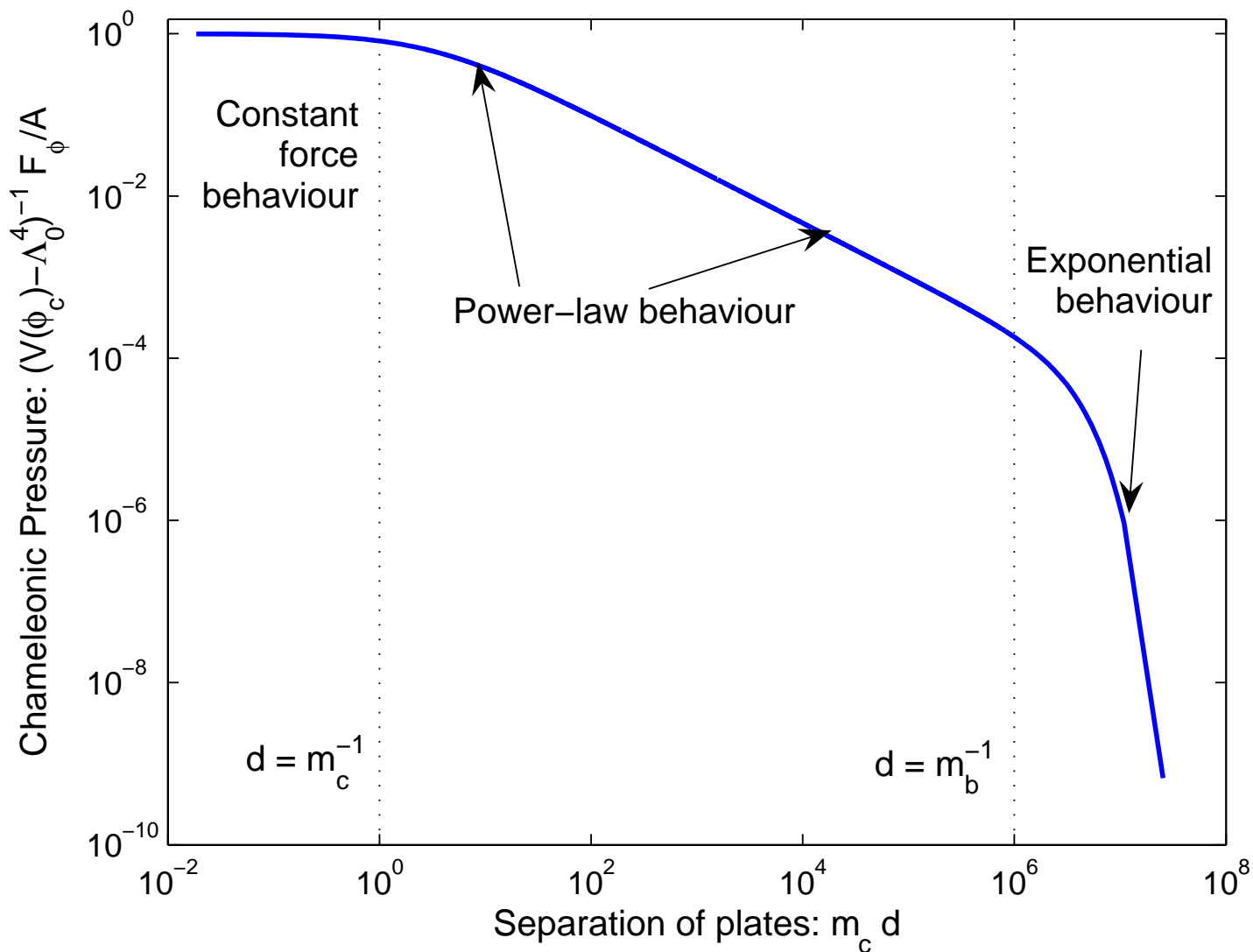
- o The force is algebraic:

$$\frac{F_\phi}{A} \sim \Lambda^4 (\Lambda d)^{-\frac{2n}{n+2}}$$

- o The dark energy scale sets a typical scale:

$$\Lambda^{-1} \sim 82 \mu m$$

Behaviour of Chameleonic Pressure for $V = \Lambda_0^4(1 + \Lambda^n/\phi^n)$; $n = 1$

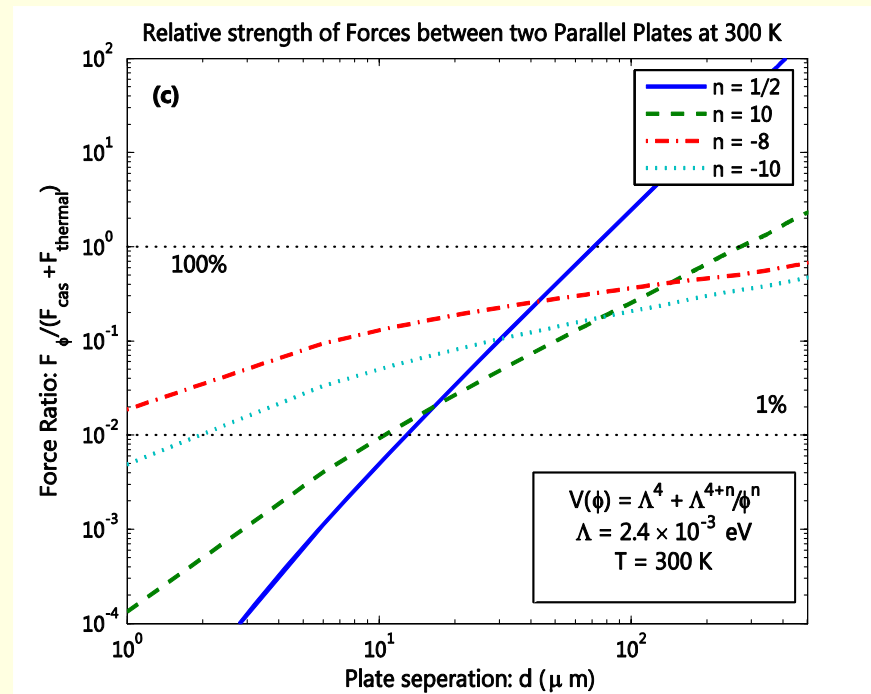


Detectability

- o The Casimir forces is also an algebraic law implying:

$$\frac{F_\phi}{F_{\text{cas}}} \sim \frac{240}{\pi^2} (\Lambda d)^{\frac{2(n+4)}{n+2}}$$

- o This can be a few percent when $d=10\mu\text{m}$ and would be 100% for $d=30\mu\text{m}$



Conclusions

- Chameleons offer the only known way of reconciling gravity tests with the existence of a dark energy particle.
- If strongly coupled to photons, they would show up in cavity experiments.
- Casimir experiments would also see a clear signal provided probed distances are large enough.

Variations of Constants

- o Scalar tensor theories predict a variation of the fundamental particle masses:

$$m_{\psi} = A(\phi)m_{\psi}^0$$

- o The gauge coupling constants are field-independent.
- o The mass of the proton is essentially a pure gluon condensation effect in Quantum ChromoDynamics, hence constant too.

- o The electron to proton ratio is field dependent:

$$\frac{m_e}{m_p} \approx A(\phi)\frac{m_e}{m_p}|_0$$

