Searching for Axion-like Particles with Active Galactic Nuclei

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Outline

- Axion-like-particles (ALPs)
- Astronomy with ALPs
- A new test for ALPs Luminosity relations
 - Results from Active Galactic Nuclei
- Conclusions

Axions and axion-like-particles

- Axions proposed by Pecci and Quinn (1977) to solve the strong CP problem of QCD
 - Pseudo-scalar particles, (approximate) shift symmetry $\phi \rightarrow \phi + \epsilon$
 - Mass proportional to coupling strength
 - Lagrangian includes $\mathcal{L} \supset \frac{\phi}{4M} \epsilon_{\mu\nu\lambda\rho} F^{\mu\nu} F^{\lambda\rho} = -\frac{\phi}{M} \mathbf{E} \cdot \mathbf{B}.$
- Axion-like-particle (ALP) is a scalar or pseudoscalar particle with a coupling to two photons
- Scalar field couples as $\mathcal{L} \supset \frac{\phi}{4M} F_{\mu\nu} F^{\mu\nu} = \frac{\phi}{2M} (\mathbf{B}^2 \mathbf{E}^2).$

- mass and coupling strength are free parameters to be constrained

Motivations and constraints

- ALPs arise in string theory and GUT models (Svrcek & Witten 2006)
- Dark energy: coupled quintessence, axionic dark energy, chameleon dark energy...

(Amendola 2000, Carroll 1998, Kim, Nilles 2003, Khoury & Weltman 2004)

• Suggested by high energy cosmic rays? See M. Fairbairn's talk

- Current experimental constraints on very light ALPs $m_{\phi} \lesssim 10^{-12} \text{ eV}.$
 - Pseudoscalar ALPs $10^{11} \text{ GeV} \leq M$ (Hagman et al. 2008)
 - Scalar ALPs $10^{26} \text{ GeV} < M$ (Will 1993)
 - Chameleonic ALPs $10^9 \text{ GeV} \lesssim M$ (CB, Davis, Shaw 2009)
 - Mass depends on the local energy density, leads to 'thin-shell' suppression of fifth forces

Optics with ALPs

• A photon travelling through a magnetic field can convert into an ALP (Raffelt & Stodolsky 1988)

$$\begin{bmatrix} \omega^2 + \partial_z^2 + \begin{pmatrix} -m_\phi^2 & \frac{B\omega}{M} & 0\\ \frac{B\omega}{M} & -\omega_P^2 & 0\\ 0 & 0 & -\omega_P^2 \end{pmatrix} \end{bmatrix} \begin{pmatrix} \phi\\ A_\perp\\ A_\parallel \end{pmatrix} = 0$$

• Probability of mixing $P = \sin^2 2\theta \sin^2 \left(\frac{\Delta(z)}{\cos 2\theta}\right)$

$$\Delta(z) = m_{eff}^2 z / 4\omega \quad \tan 2\theta = 2B\omega / M m_{eff}^2 \quad m_5^{\text{ell}} = |m_5^{\phi} - m_5^2|$$

- Increases with frequency
- At large frequencies approaches an upper bound and becomes frequency independent
- So far no detection of ALPs in the laboratory

Astronomy with ALPs

- Can we use astronomy to search for ALPs?
- Magnetic fields exist in galaxies and galaxy clusters
 - Disadvantages: Field fluctuates on many different scales
 Fields typically weaker than in the lab
 - Advantages: Fields extend over much greater distances
- Sufficient to use the cell magnetic field model



- Strong mixing occurs when $NP \gg 1$ $N\Delta(L) \leq \pi/2$
 - Probability of mixing becomes large and frequency independent
 - After passing through a large number of domains the initial flux becomes, on average, equally distributed between A_1 , A_2 ϕ
 - The attenuation factor $C = S_0(N)/S_0(0)$

Strong mixing

$$f_{C}(c;p_{0}) = \frac{1}{\sqrt{1-p_{0}^{2}}} \left[\tan^{-1} \left(\sqrt{a} \left(1 - \frac{2c_{+}}{1+p_{0}} \right)^{-1/2} \right) \right]$$

$$-\tan^{-1} \left(\sqrt{a} \left(1 - \frac{2c_{1}}{1-p_{0}} \right)^{1/2} \right) \right]$$

$$a = (1+p_{0})/(1-p_{0})$$

$$c_{\pm} = \min (c, (1 \pm p_{0})/2)$$

$$\int_{0}^{\infty} \frac{1}{2} \int_{0}^{0} \frac{1}{1-p_{0}} \int_{0}^{0}$$

С

A new test for ALPs - Luminosity relations

- Problem: We do not generally know the expected high energy flux from astronomical sources
- For a number of types of object there exist 'luminosity relations' correlating their high and low frequency luminosity
 - At low frequencies light mixes weakly with ALPs, and the luminosity is not altered at leading order
- Relations take the form $\log_{10} Y_i = a + b \log_{10} X_i + S_i$ High energy flux low energy flux scatter - Standard to model the scatter as normally distributed

$$S_i = \sigma \delta_i \qquad \delta \sim N(0, 1)$$

- If high frequency light mixes with ALPs $S_i = \sigma \delta_i - \log_{10} C_i$

A new test for ALPs - Luminosity relations

• Compare the two models with the Likelihood ratio test

$$r(p_0) = 2\log\left(\frac{\hat{L}_1(p_0)}{\hat{L}_0}\right)$$

- We consider photons propagating through the magnetic fields of galaxy clusters $B \approx 1 - 10 \mu \text{G}$ $L \sim \text{kpc}$ $N \approx 100 - 1000$
- If $M \lesssim 10^{11} \text{ GeV}$ strong mixing for x-ray and gamma-ray photons

- Assume
$$m_{\phi}^2 \lesssim \omega_F^2$$

 $\omega_{\rm pl} \sim 10^{-12} \, {\rm eV}$

Active galactic nuclei

- Luminosity relation between 2 keV x-ray luminosity and 5 eV optical luminosity
 - Observations of 77 AGN from COMBO-17 and ROSAT surveys and 126 objects from the SDSS survey (Steffen et al. 2006,Strateve 2005)
- For the combined data set $r(p_0 \leq 0.5) \approx 25$
 - expectation from AGN physics is $p_0 < 0.1$
- Qualitative check independent of any null hypothesis
 - Construct 10⁵ new data sets (with 203 points) by bootstrap resampling with replacement
- Compute statistical moments for data sets k_m
 - k_2 the rms mean, k_3^3/k_2^3 the skew, ...

Active galactic nuclei - Fingerprints



Active galactic nuclei - Fingerprints



Conclusions & future tests

- AGN luminosity relations show very strong evidence for ALP-photon mixing over null hypothesis of Gaussian noise
- There is structure in the scatter which, qualitatively, looks like strong ALP-photon mixing $M \lesssim 10^{11}$ GeV
 - pseudoscalar, or chameleonic field
- $m_{\phi} \lesssim 10^{-12} \text{ eV}.$
- Qualitative similarity also exists for higher moments of the distribution
- We do not understand the physics of the luminosity relation this could be standard physics mimicking ALPs
 - No evidence for correlation between redshift and scatter, so cannot be:
 - evolution effects
 - inaccurate choice of cosmological model

Conclusions & future tests

- Is this effect independent of the objects studied?
 - Luminosity relations also exist currently for GRBs and Blazars
 - Likelihood ratio test gives $r \approx 1.6$
 - either too few data points
 - or too large intrinsic scatter
- If the explanation is strong mixing expect to also induce a large linear polarization
 - Could be measured with e.g. proposed International X-ray Observatory
- CAST should be sensitive to these couplings in the next decade

Schaefer (2007)

Bloom, (2007)

Xie, Zhang, Fan (1997)



Strong mixing

- Strong mixing occurs when $NP \gg 1$ $N\Delta(L) \lesssim \pi/2$ - Probability of mixing becomes large and frequency independent
- Mixing between ALP and photons, and between different components of the photon becomes large
 - After passing through a large number of domains the initial flux becomes, on average, equally distributed between A_1 , A_2 ϕ
- If photon component fully polarized, initially $\mathbf{u}(0) = (\phi(0), A_1(0), A_2(0))^T |\mathbf{u}(0)| = 1$
- Finally

$$\mathbf{u}_N = (x, \sqrt{1 - x^2} \cos \pi \Theta, \sqrt{1 - x^2} \sin \pi \Theta)^T \quad x, \Theta \sim U(-1, 1)$$

- But in astronomy our sources are partially polarized or unpolarized

Strong mixing

• To describe a partially or unpolarized light beam use the Stokes vector $(S_0) = (|A_1|^2 + |A_2|^2)$

$$\mathbf{S} \equiv \begin{pmatrix} S_0 \\ S_1 \\ S_2 \\ S_3 \end{pmatrix} = \begin{pmatrix} |A_1|^2 + |A_2|^2 \\ |A_1|^2 - |A_2|^2 \\ A_1\bar{A}_2 + \bar{A}_1A_2 \\ i(A_1\bar{A}_2 - \bar{A}_1A_2) \end{pmatrix}$$

- Initial photon state $\mathbf{S}(0) = \frac{1}{2}(1+p_0) \begin{pmatrix} 1\\ 1\\ 0\\ 0 \end{pmatrix} + \frac{1}{2}(1-p_0) \begin{pmatrix} 1\\ -1\\ 0\\ 0 \end{pmatrix}$
- Final photon state

$$\mathbf{S}(N) = \frac{1}{2}(1+p_0)(1-x^2) \begin{pmatrix} 1\\\cos 2\theta\\\sin 2\theta\\0 \end{pmatrix} + \frac{1}{2}(1-p_0)(1-y^2) \begin{pmatrix} 1\\\cos 2\theta\\\sin 2\theta\\0 \end{pmatrix}$$

A new test for ALPs - Luminosity relations

• Given data set $\{X_i, Y_i\}$ the likelihood of the models is

$$L_f(a, b, \sigma; p_0) = \prod_i \frac{1}{\sqrt{2\pi}} \sigma \int_0^1 e^{-\frac{z_i^2}{2\sigma^2}} f_C(c; p_0) \, dc$$

$$z_i = \log_{10} Y_i - a - b \log_{10} X_i - \log_{10} ((1 - f) + fc)$$

-f = 1 when high frequency light is strongly mixed with ALPs -f = 0 when high frequency light does not mix with ALPs

- For each model fit for a, b, σ by maximising the likelihood
- Define the ratio $r(p_0) = 2 \log \left(\frac{\hat{L}_1(p_0)}{\hat{L}_0}\right)$

- \hat{L} the model with the most likely parameters

Active galactic nuclei

- Results only weakly dependent on initial polarization,
 - robust to relaxing the universality of p_0
- In the ALP strong mixing model we have taken the probability of mixing $P_{\text{mix}} = 1$ to be universal
 - result robust to different values of $P_{\rm mix}$
 - Data only constrains $P_{\text{mix}} \ge 0.08$ at 95% confidence
- Have we just chosen a bad null hypothesis?
 - Null hypothesis was chosen because of its simplicity and because for similar relations it is a good fit
 - Do not understand the cause of the luminosity relation for AGN, so we cannot improve the null hypothesis

Discussion

- Qualitative similarity also exists for higher moments of the distribution
- We do not understand the physics of the luminosity relation this could be standard physics mimicking ALPs
 - No evidence for correlation between redshift and scatter, so cannot be:
 - evolution effects
 - inaccurate choice of cosmological model
- Luminosity relations also exist for GRBs and Blazars
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Conclusions

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- pseudoscalar, or chameleonic field
- But because we do not understand AGN physics cannot rule out explanation in terms of old physics
- If the explanation is strong mixing expect to also induce a large linear polarization
 - Could be measured with e.g. proposed International X-ray Observatory
- If we are seeing a pseudoscalar ALP hopefully seen at CAST in the next decade