

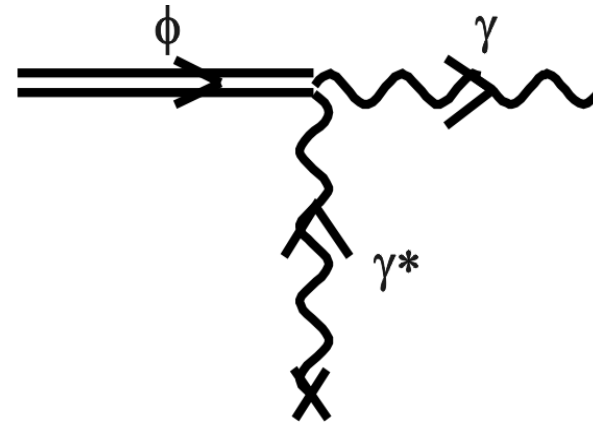
Search for Solar ALPs in the Low Energy Range at CAST



Outline:

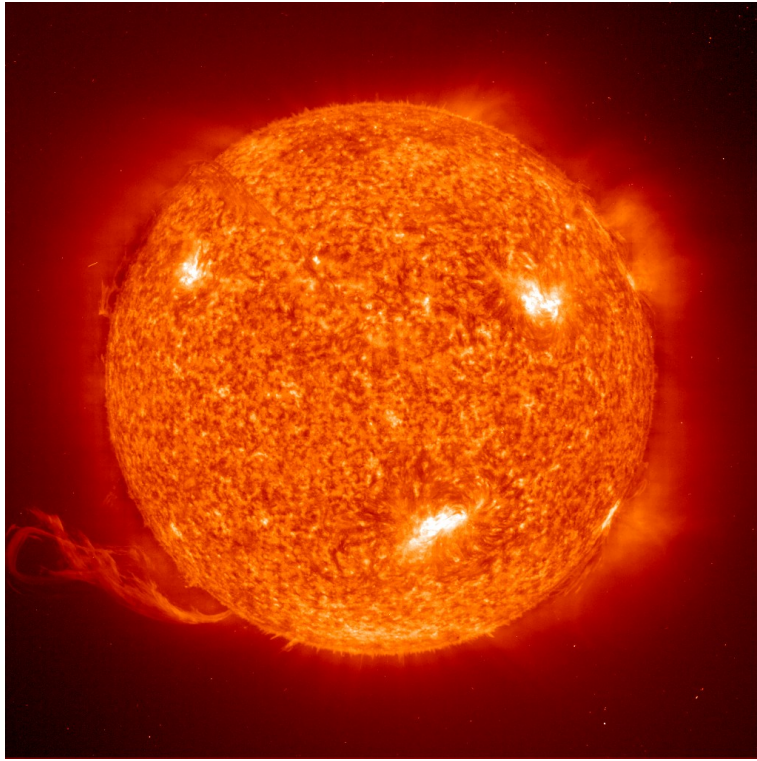
- Introduction**
 - ALPs**
 - Sun**
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- BaRBE**
- Experimental setup**
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- Summary**

Introduction: ALPs



- couple to photons via Primakoff effect
 - nearly massless
 - weakly interacting
- The axion is a hypothetical elementary particle postulated by the Peccei-Quinn theory in 1977 to resolve the strong-CP problem.
- In particle physics, the strong CP problem is the puzzling question why quantum chromodynamics (QCD) does not seem to break the CP symmetry.

Introduction: Sun



Volume: $1.412 \times 10^{18} \text{ km}^3$ _ 1,300,000 Earth's
Mass : $1.989 \times 10^{30} \text{ kg}$ _ 332,900 x Earth's
Radius : $6.955 \times 10^5 \text{ km}$ _ 109 x Earth's
Mean Distance to Earth: 149.60 million km (1AU)
Solar Constant (Total Solar Irradiance): 1.365 - 1.369 kW/m²
 (at the mean distance of the earth from the Sun, about one AU)
- solar corona heating

*AXIONS: RECENT SEARCHES AND NEW LIMITS, GEORG G. RAFFELT, arXiv:hep-ph/0504152 v2 2 May 2005

Integrating over a standard solar model, one finds an axion flux at Earth that is well approximated by (E in keV)

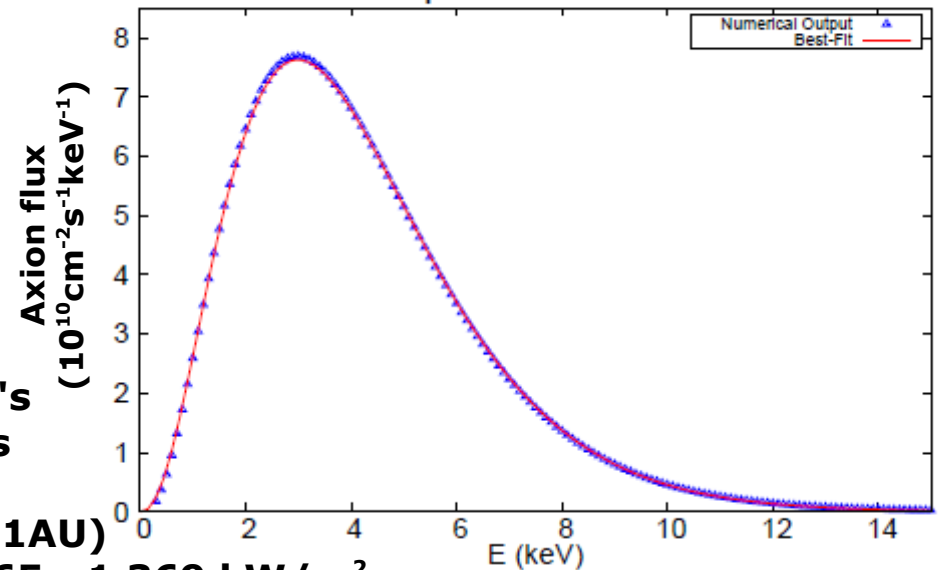
$$\frac{d\Phi_a}{dE} = g_{10}^2 6.0 \times 10^{10} \text{ cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1} E^{2.481} e^{-E/1.205}, \quad (13)$$

where $g_{10} = g_{a\gamma\gamma}/(10^{-10} \text{ GeV}^{-1})$. The integrated flux parameters are

$$\begin{aligned} \Phi_a &= g_{10}^2 3.75 \times 10^{11} \text{ cm}^{-2} \text{ s}^{-1}, \\ L_a &= g_{10}^2 1.85 \times 10^{-3} L_{\odot}. \end{aligned} \quad (14)$$

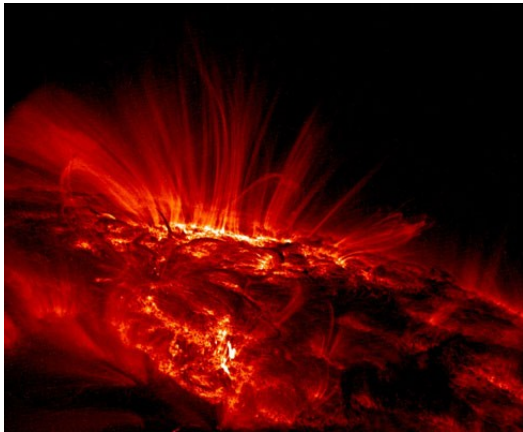
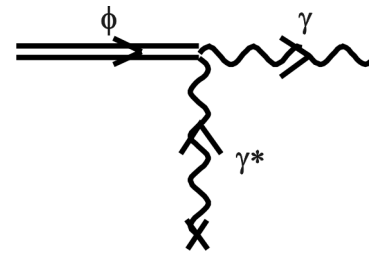
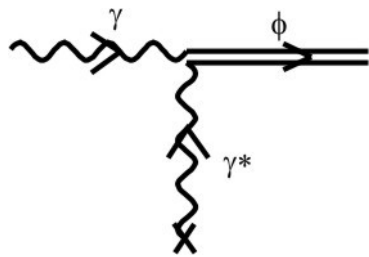
The maximum of the distribution is at 3.0 keV, the average energy is 4.2 keV.

Solar Axion Spectra - 2004 Solar Model

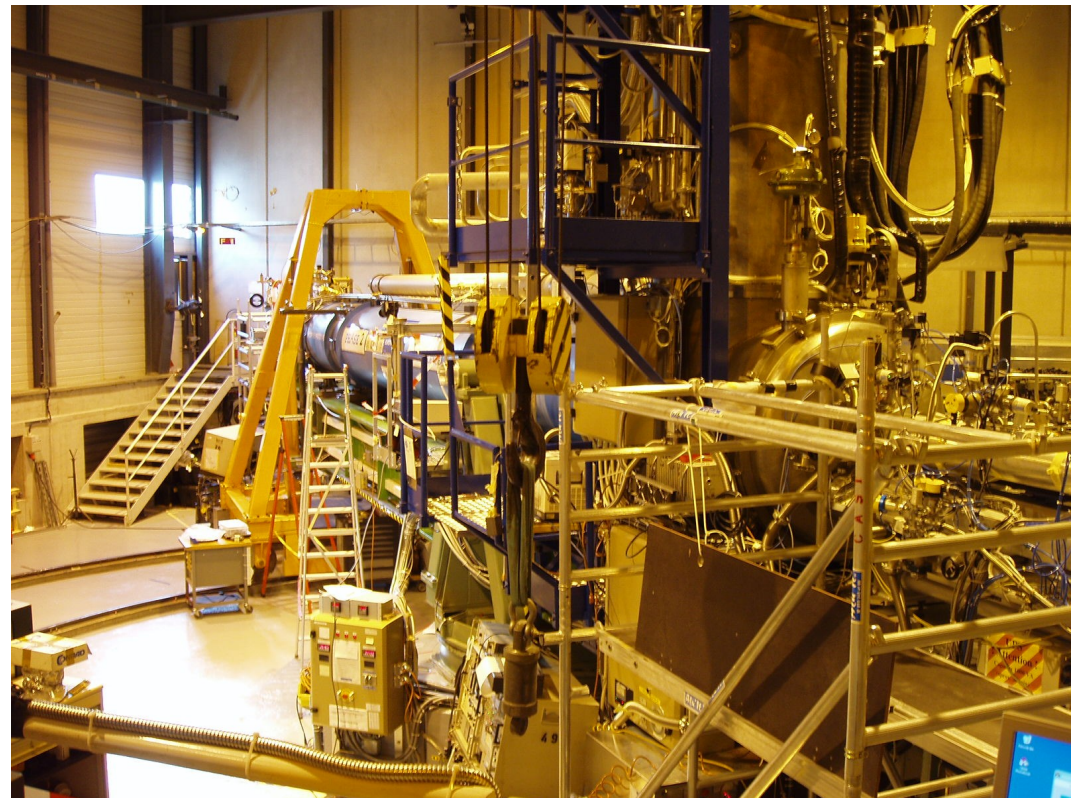


$$\times \left(\frac{g_{a\gamma\gamma}}{10^{-10} \text{ GeV}^{-1}} \right)^2$$

Introduction: CAST



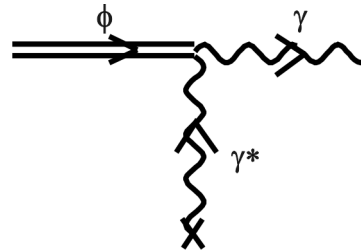
- reconverted photons peaked at 3.2 keV
- LHC dipole magnet $L=9.3$ m, $B=9$ T
 $d=43$ mm, $A=14.5$ cm² x2
- data taking 1.5 h at sunset & sunrise
- 3 MM & 1 CCD
- PMT to look for low energy reconverted photons



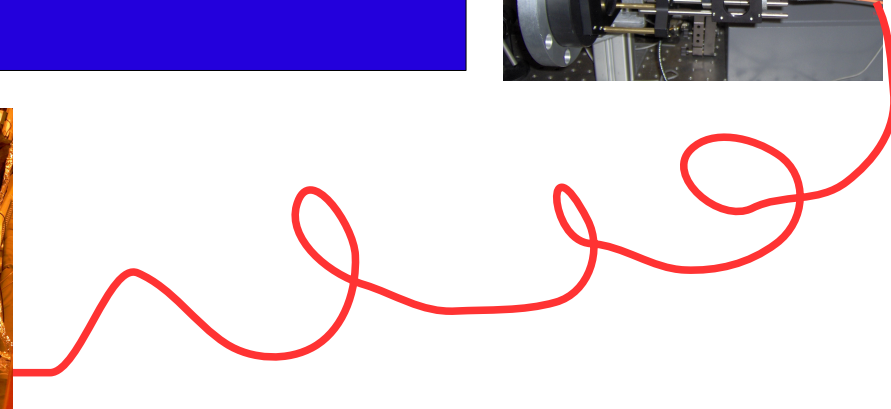
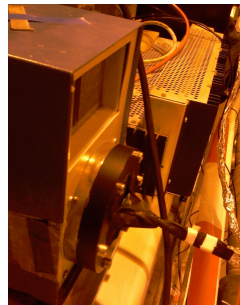
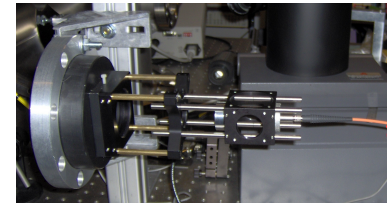
- view from BaRBE counting room

BaRBE

Low Rate – Low Energy



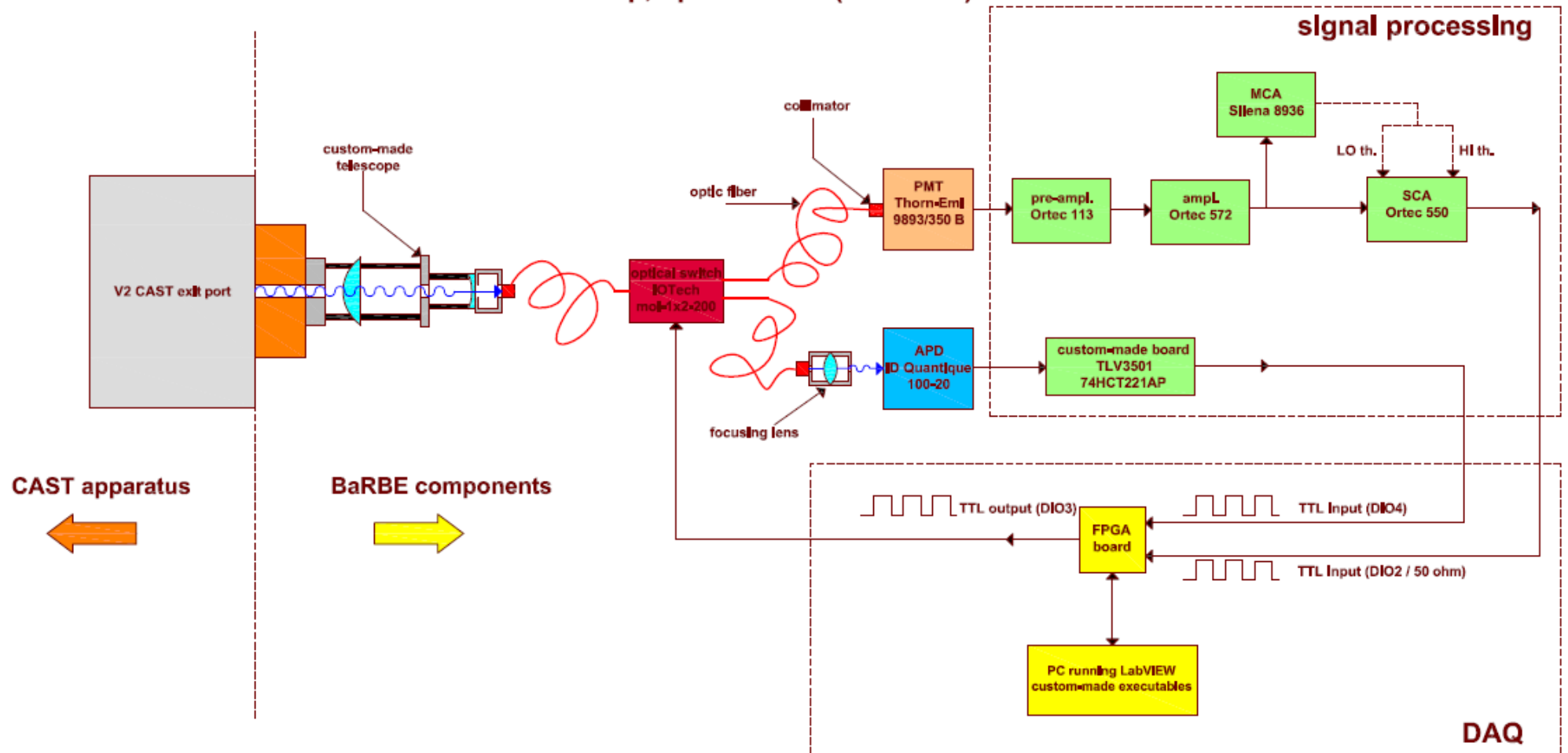
MAGNET



- limited experimental effort financed by INFN
- study various types of detectors with low background
- check coupling with the reconversion magnet with an optical fibre
- look for low energy photons in the magnet bore
- first tests with PMT with an ALPs beam

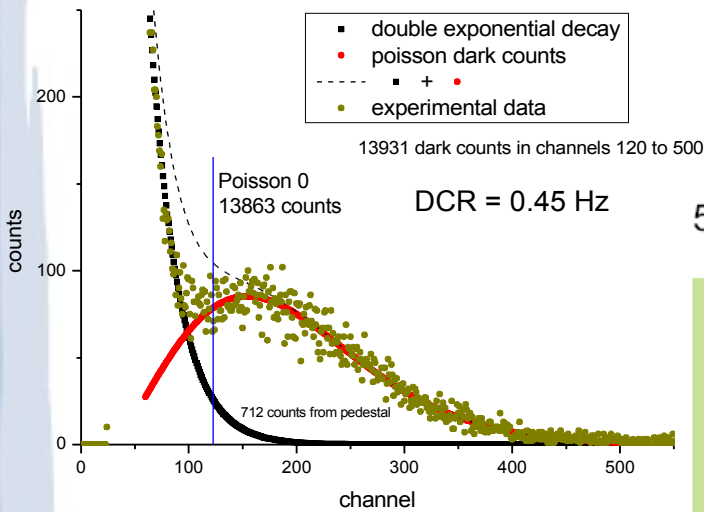
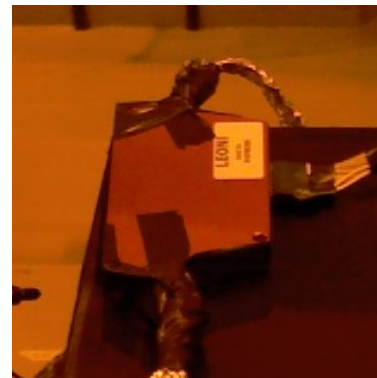
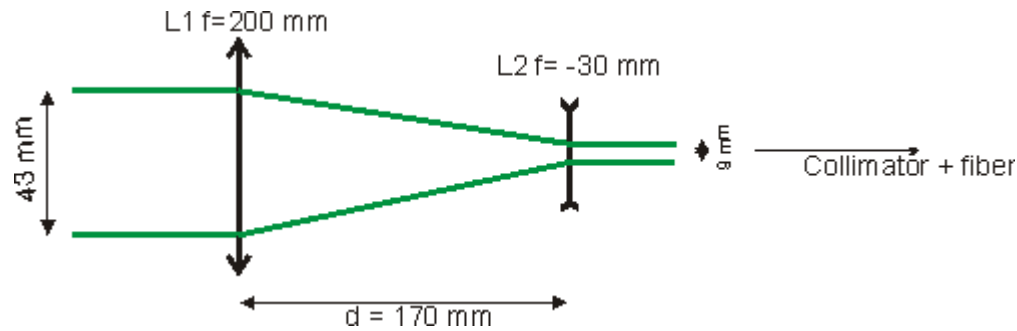
Experimental setup

BaRBE measurement set-up, option No.1 (Oct. 2007)

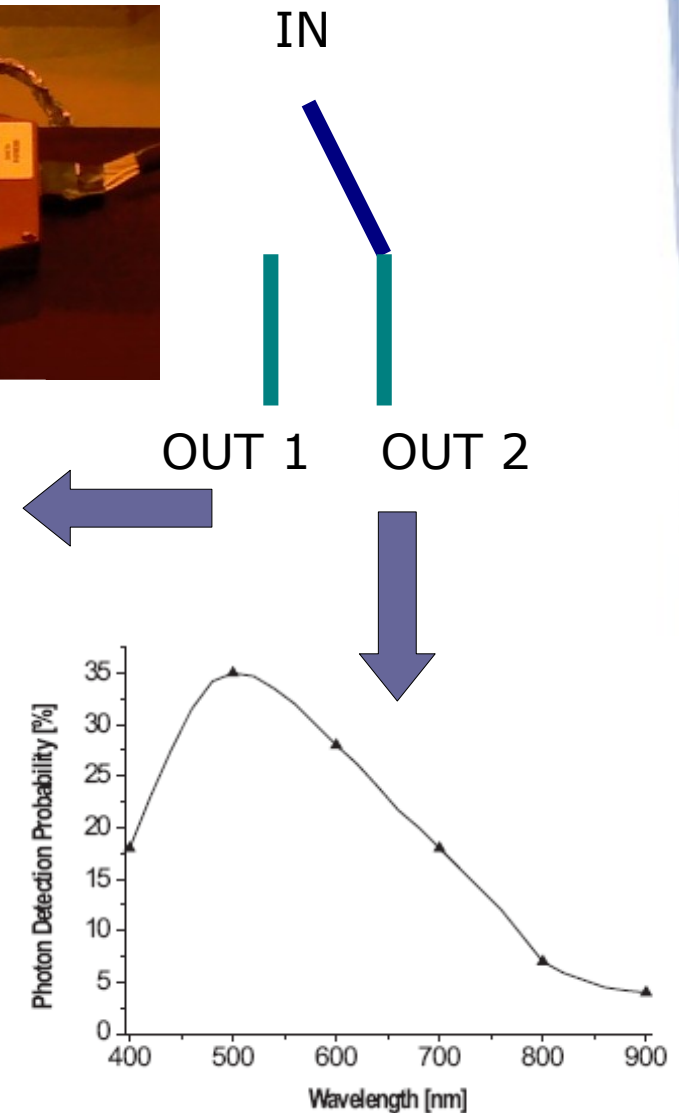
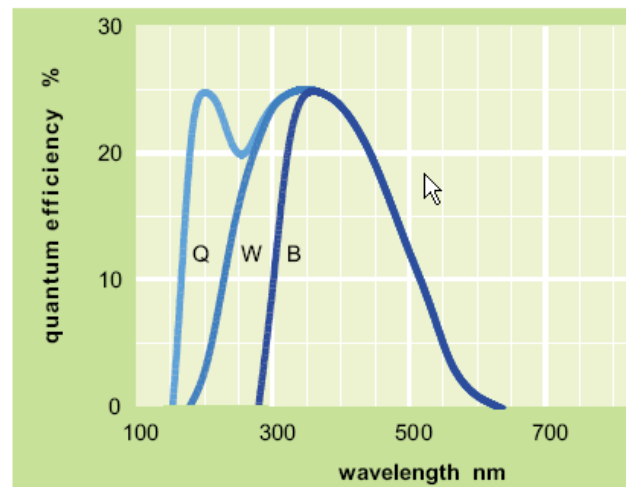


- main idea: divide beam time between two detectors
- use an optical switch
- simultaneous background/live data taking

Experimental setup

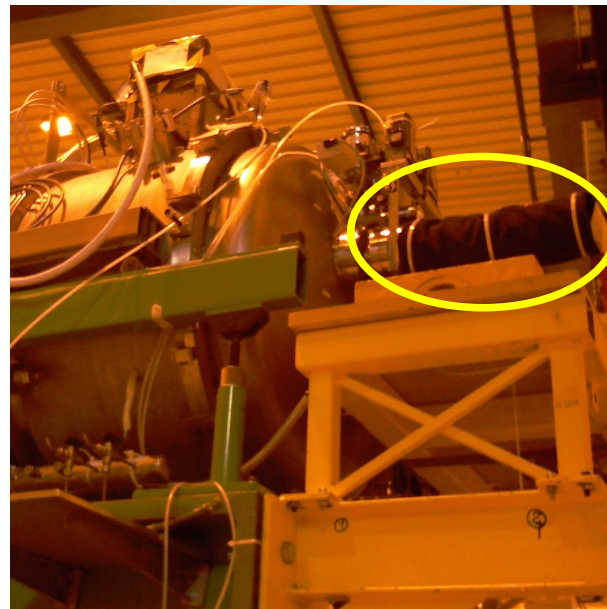
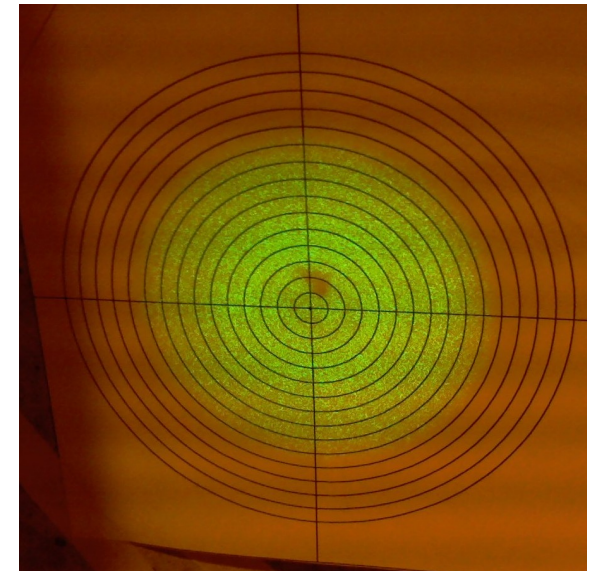
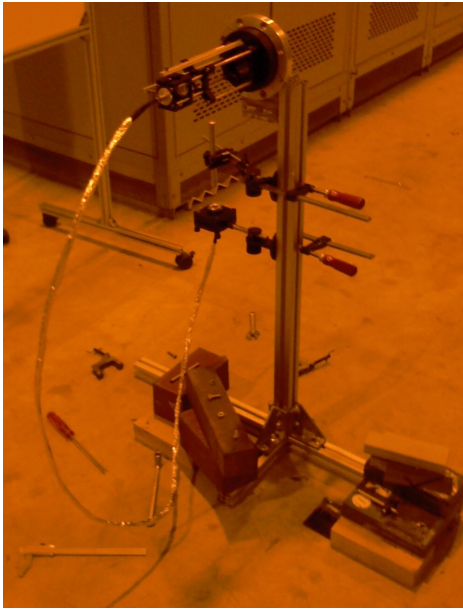


5 typical spectral response curve

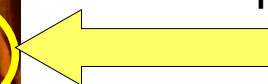


- PMT Thorn EMI 9893/350B, measured dark count rate 0.4 Hz
- APD idQuantique model id100-20, 0.4 Hz DCR, 20 micron diameter

Experimental setup: Alignment



Telescope on CAST magnet



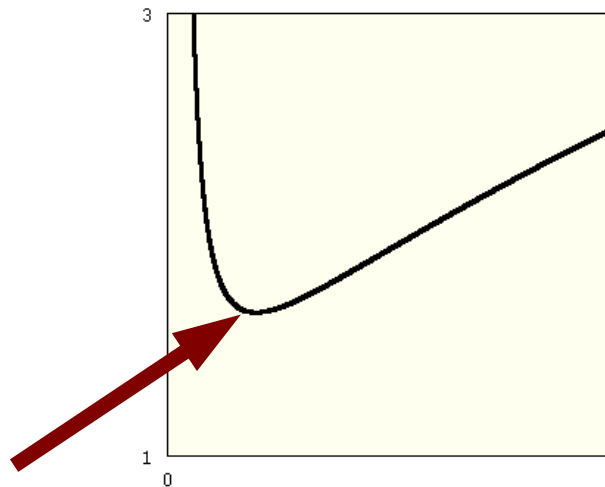
- blind alignment, with surveyors help
- take reference points on the telescope with respect to the beam, align the reference points with the magnet axis

Analysis

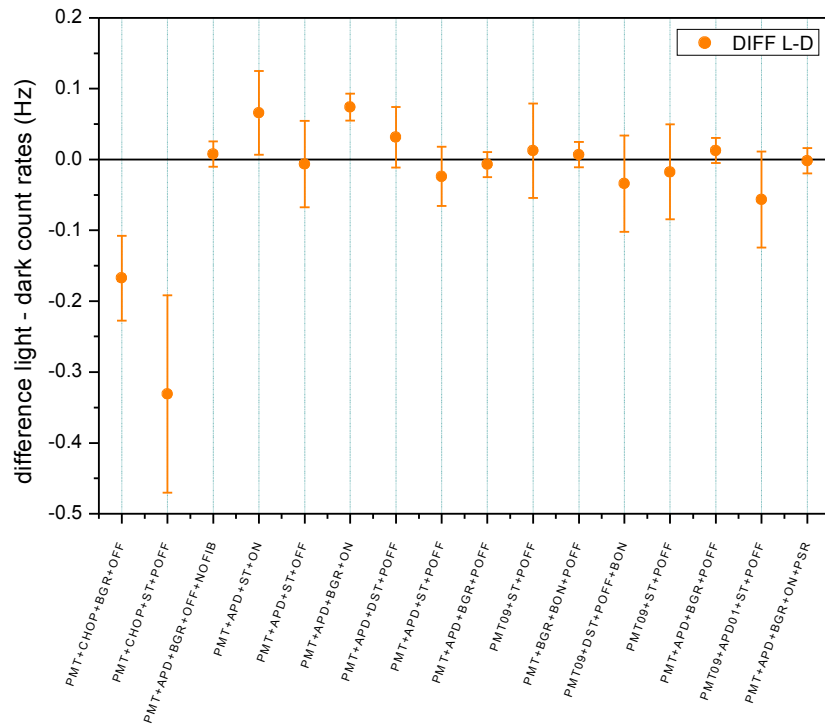
Since the data follow a Poisson distribution, another way to calculate the mean rate of counts (light and dark) is to solve the following equation (for $x = 0$): $N_x = A \cdot e^{(-m)} m^x / x!$ where x is the channel number, A is the total number of occurrences in all channels, N_x is the number of occurrences in the x - th channel measured experimentally. The advantage of this method is that occurrences in the channel 0 are not affected by the afterpulses so the obtained value does not have to be corrected for.

The obtained rate is afterwards normalized to 1 s period and is given in Hz.

The error is affected by the choice of A (function of switching frequency and acquisition time) and is given by $dm_n = \sqrt{(A)} \cdot e^{(N/2A)} / T_D$. This function has a minimum for $A = N$, i.e. the error is minimized when the number of counting periods is equal to the number of photons N counted during the acquisition time T_D as shown in the following graph ($A = N = 1$).



Results



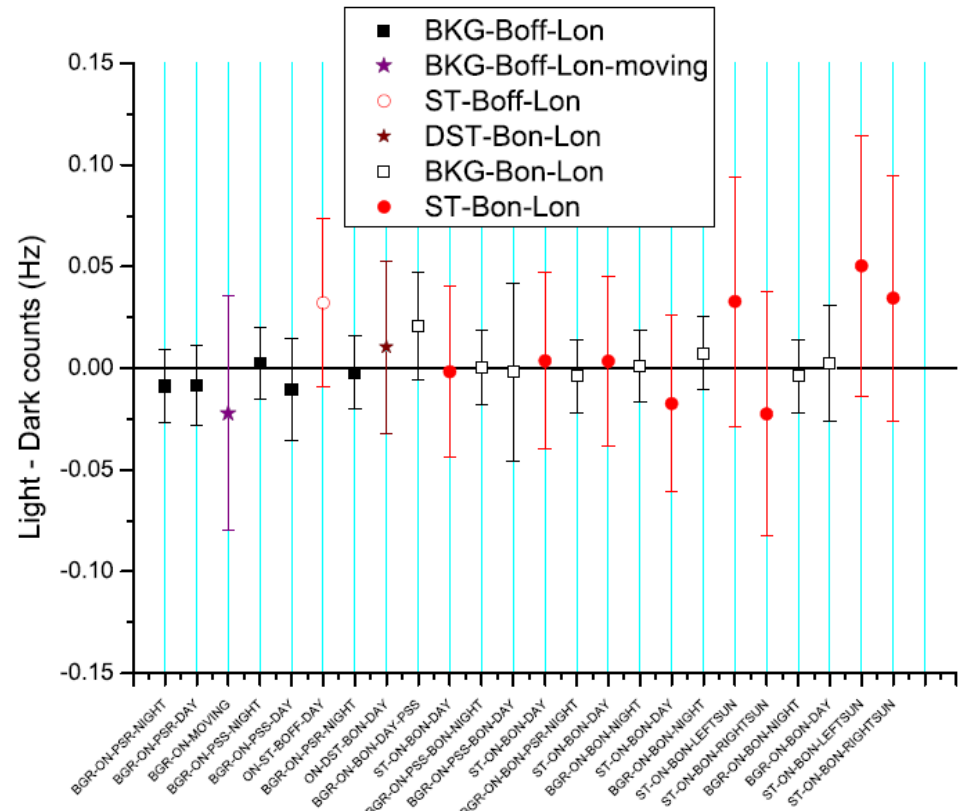
November 2007

$$\tau = 50000 \text{ s}$$

$$g_{a\gamma\gamma} \sim 10^{-6} \text{ GeV}^{-1}$$

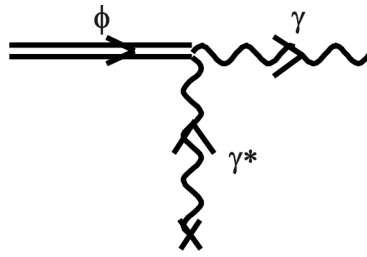
$$\text{DCR} = 0.35 \text{ Hz} \pm 0.02 \text{ Hz}$$

March 2008

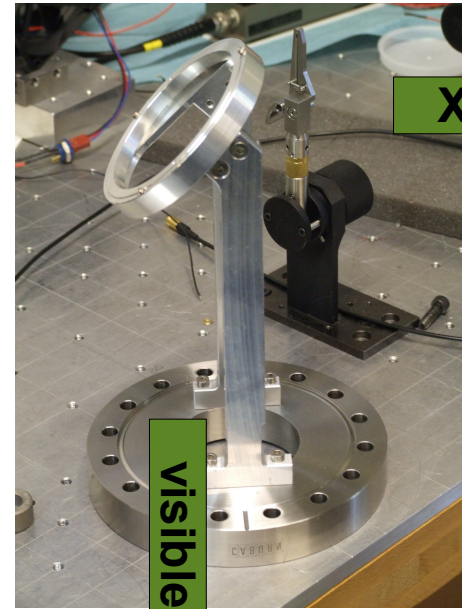


- there is no signal in both data taking runs
- the background is the same as in laboratory conditions

But ...

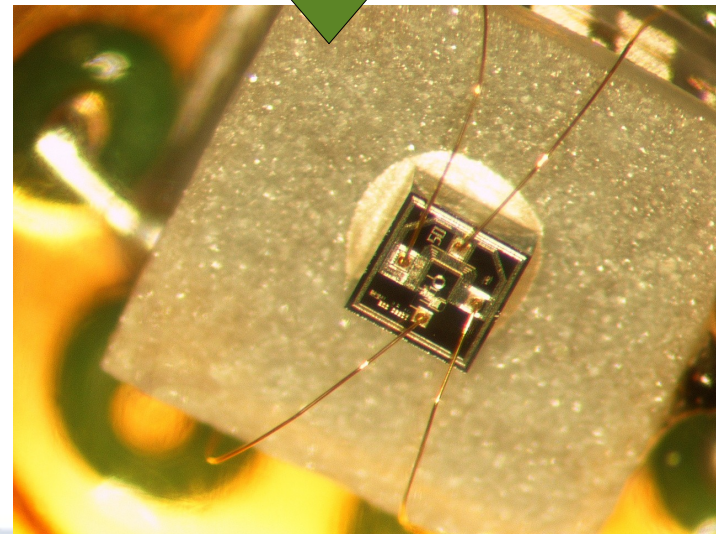


- ALPs source – black box
- semitransparent mirror
- continue to take data
- lower DCR --> detector choice

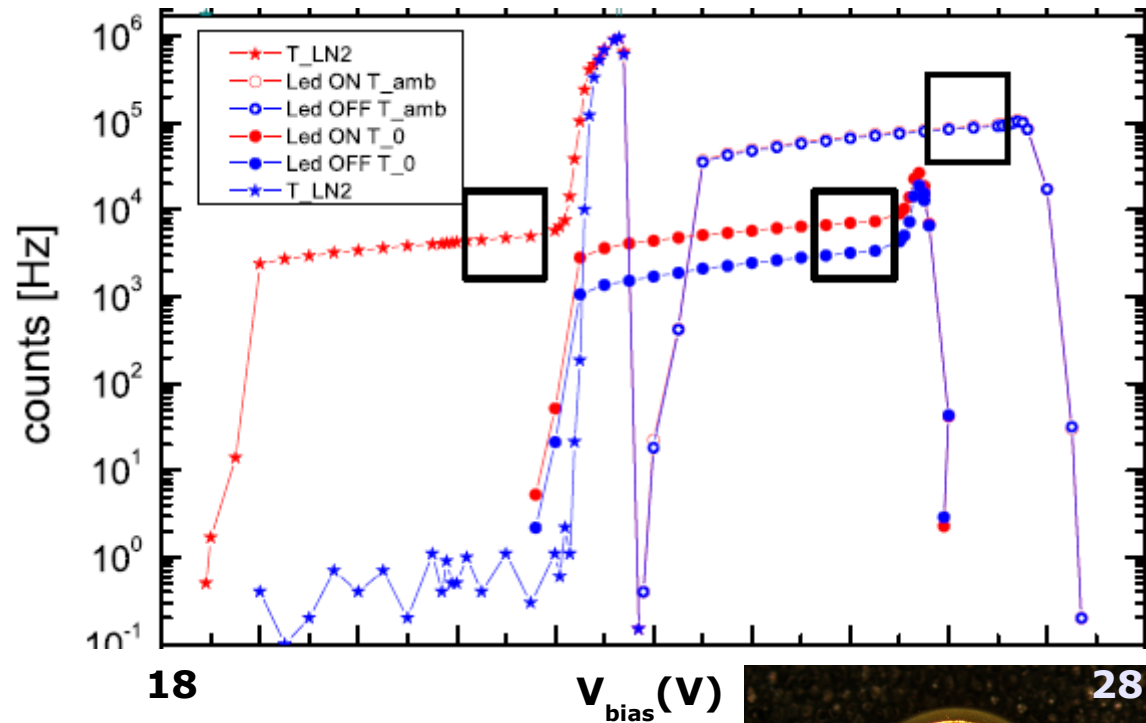
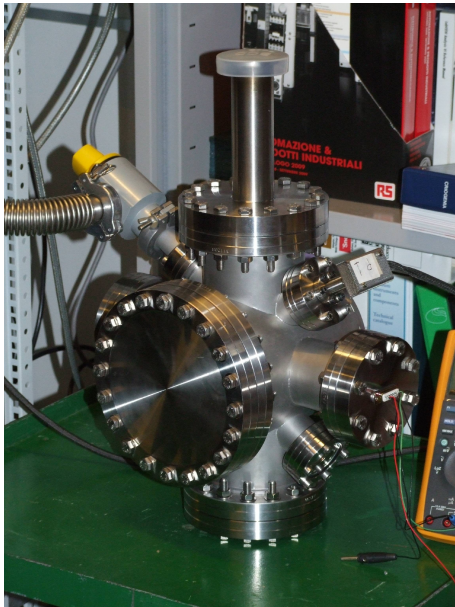


X rays →

↓ visible



Poor (wise) man's choice: Cooled APD



- 10^4 background reduction
- no efficiency loss
- works for "bad" APDs
- will it work for "good" ones (starting DCR 1 Hz) ?

Summary

- 10^4 background reduction with cooled APD
- experimental setup is working
- with semitransparent mirror it is possible to continue taking data at CAST
- resonant regeneration
- a surprise can exit from the ALPs black box

